

# Morse indices and the number of blow up points of blowing-up solutions for a Liouville equation with singular data

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**Abstract.** Let  $\Omega \subset \mathbb{R}^2$  be a smooth bounded domain and let  $\Gamma = \{p_1, \dots, p_N\} \subset \Omega$  be the set of prescribed points. Consider the Liouville type equation

$$-\Delta u = \lambda \prod_{j=1}^N |x - p_j|^{2\alpha_j} V(x) e^u \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega,$$

where  $\alpha_j$  ( $j = 1, \dots, N$ ) are positive numbers,  $V(x) > 0$  is a given smooth function on  $\overline{\Omega}$ , and  $\lambda > 0$  is a parameter. Let  $\{u_n\}$  be a blowing up solution sequence for  $\lambda = \lambda_n \downarrow 0$  having the  $m$ -points blow up set  $S = \{q_1, \dots, q_m\} \subset \Omega$ , i.e.,

$$\lambda_n \prod_{j=1}^N |x - p_j|^{2\alpha_j} V(x) e^{u_n} dx \rightharpoonup \sum_{i=1}^m b_i \delta_{q_i}$$

in the sense of measures, where  $b_i = 8\pi$  if  $q_i \notin \Gamma$ ,  $b_i = 8\pi(1 + \alpha_j)$  if  $q_i = p_j$  for some  $p_j \in \Gamma$ . We show that the number of blow up points  $m$  is less than or equal to the Morse index of  $u_n$  for  $n$  sufficiently large, provided  $\alpha_j \in (0, +\infty) \setminus \mathbb{N}$  for all  $j = 1, \dots, N$ . This is a generalization of the result [14] in which nonsingular case ( $\alpha_j = 0$  for all  $j$ ) was studied.

**Keywords:** Liouville equation, blow up points, singular data, concentration compactness result, Morse indices.

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## 1 Introduction

Let  $\Omega$  be a smooth bounded domain in  $\mathbb{R}^2$  and  $\lambda > 0$  is a parameter. Motivated by some physical problems in selfdual Gauge Field Theories such as

Chern-Simons vortex theories or others (see [12], [15]), some researchers are interested in the analysis of the problem

$$\begin{cases} -\Delta v = \lambda e^v - 4\pi \sum_{j=1}^N \alpha_j \delta_{p_j} & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega, \end{cases} \quad (1.1)$$

where  $\Gamma = \{p_1, \dots, p_N\} \subset \Omega$  is the set of prescribed singular sources (called “vortices”),  $\delta_p$  is a Dirac mass supported at  $p$ , and  $\alpha_j > 0$ .

If we introduce the Green’s function of  $-\Delta$  acting on  $H_0^1(\Omega)$ :

$$\begin{cases} -\Delta_x G(x, p) = \delta_p & \text{for } x \in \Omega, \\ G(x, p) = 0 & \text{for } x \in \partial\Omega, \end{cases}$$

and write  $G(x, p) = \frac{1}{2\pi} \log |x - p|^{-1} + H(x, p)$ , where  $H(x, p)$  is the regular part of  $G$ , then the problem (1.1) is equivalent to

$$\begin{cases} -\Delta u = \lambda \prod_{j=1}^N |x - p_j|^{2\alpha_j} V(x) e^u & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega \end{cases} \quad (1.2)$$

where  $u = v + 4\pi \sum_{j=1}^N \alpha_j G(x, p_j)$  and  $V(x) = e^{-4\pi \sum_{j=1}^N \alpha_j H(x, p_j)}$  is a smooth positive function on  $\bar{\Omega}$ . By this reason, we are led to consider the problem (1.2) for general smooth positive functions  $V$ . In this case, the study of asymptotic behavior of solutions  $u_n$  for  $\lambda = \lambda_n \rightarrow +0$  in (1.2) was done by P. Esposito in [5] (see also [6] [7]), which extends the results of [9], [10] where the regular case ( $\alpha_j = 0, \forall j$ ) was considered.

**Theorem 1** (*P. Esposito*) *Let  $V$  be a smooth positive function on  $\bar{\Omega}$  and set  $K(x) = \prod_{j=1}^N |x - p_j|^{2\alpha_j} V(x)$ . Let  $\{\lambda_n\}$  be a sequence of positive numbers with  $\lambda_n \rightarrow 0$  and let  $\{u_n\}$  be a solution sequence of (1.2) for  $\lambda = \lambda_n$  such that*

$$\Sigma_n = \lambda_n \int_{\Omega} K(x) e^{u_n} dx = O(1) \quad \text{as } n \rightarrow \infty.$$

*Then the following alternative holds:*

- (i) *If  $\Sigma_n \rightarrow 0$  as  $n \rightarrow \infty$ , then  $u_n \rightarrow 0$  in  $C^{2,\alpha}(\Omega)$  for some  $\alpha \in (0, 1)$  and  $u_n$  coincides with the unique minimal solution of (1.2).*

(ii) If  $\Sigma_n \rightarrow L$  for some  $L \neq 0$ , then (up to subsequence) there exists a nonempty finite set  $S = \{q_1, \dots, q_m\} \subset \Omega$  (blow up set) such that  $\{u_n\}$  is uniformly bounded in  $L_{loc}^\infty(\bar{\Omega} \setminus S)$ , and

$$\lambda_n K(x) e^{u_n} dx \rightharpoonup \sum_{i=1}^m b_i \delta_{q_i} \quad \text{in the sense of measures,} \quad (1.3)$$

$$u_n \rightarrow \sum_{i=1}^m b_i G(\cdot, q_i) \quad \text{in } C_{loc}^2(\bar{\Omega} \setminus S) \quad (1.4)$$

as  $n \rightarrow \infty$ , where  $b_i = 8\pi$  if  $q_i \notin \Gamma$ ,  $b_i = 8\pi(1 + \alpha_j)$  if  $q_i = p_j$  for some  $p_j \in \Gamma$ .

Furthermore, as for the location of blow up points in the case (ii), we have the following:

If  $S \cap \Gamma = \emptyset$ , then  $(q_1, \dots, q_m)$  is a critical point for the function

$$\mathcal{F}(x_1, \dots, x_m) = \sum_{i=1}^m H(x_i, x_i) + \sum_{i,j=1, i \neq j}^m G(x_i, x_j) + \frac{1}{4\pi} \sum_{i=1}^m \log K(x_i).$$

If  $S \cap \Gamma = \{p_{j_1}, \dots, p_{j_s}\}$  and  $S \setminus \Gamma = \{q_{i_1}, \dots, q_{i_k}\}$  with  $s + k = m$ , then  $(q_{i_1}, \dots, q_{i_k})$  is a critical point for the function

$$\tilde{\mathcal{F}}(x_1, \dots, x_k) = \mathcal{F}(x_1, \dots, x_k) + \mathcal{G}(x_1, \dots, x_k; p_{j_1}, \dots, p_{j_s}),$$

where

$$\mathcal{G}(x_1, \dots, x_k; a_1, \dots, a_s) = \frac{1}{4\pi} \left( \sum_{i=1}^k \sum_{j=1}^s 8\pi(1 + \alpha_j) G(x_i, a_j) \right).$$

Also, as a vice versa of Theorem 1, Esposito constructed blowing up solutions with a prescribed blow up set  $S$  under the additional assumption that  $\alpha_j \in (0, +\infty) \setminus \mathbb{N}$  for all  $j = 1, \dots, N$ ; see [6].

In the following, let  $i_M(u)$  denote the Morse index of a solution  $u$  of (1.2), i.e., the number of negative eigenvalues of the operator  $L_u = -\Delta - \lambda K(x) e^u$  acting on  $H_0^1(\Omega)$ .

Now, we state the main result of this note, which is a generalization of [13] [14] in this case.

**Theorem 2** *Let  $\{u_n\}$  be a solution sequence of (1.2) for  $\lambda = \lambda_n$  with  $\Sigma_n = O(1)$  as  $n \rightarrow \infty$  and let  $S = \{q_1, \dots, q_m\}$  be its blow up set (possibly  $S = \emptyset$ ). Assume  $\alpha_j \in (0, +\infty) \setminus \mathbb{N}$  for all  $j = 1, \dots, N$ . Then  $m \leq i_M(u_n)$  for  $n$  sufficiently large.*

As a corollary, we obtain the following assertion.

**Corollary 3** *Let  $\{u_n\}$  be a solution sequence of (1.2) for  $\lambda = \lambda_n$  with  $\Sigma_n = O(1)$  as  $n \rightarrow \infty$ . Assume  $\alpha_j \in (0, +\infty) \setminus \mathbb{N}$  for all  $j = 1, \dots, N$  and the Morse index  $i_M(u_n) = 1$  for any  $n$  large. Then the number of blow up points of  $\{u_n\}$  is exactly 1.*

*Proof.* By Theorem 2 and the assumption that  $i_M(u_n) = 1$  for  $n$  large, we see that the number of blow up points  $\#S$  is 0 or 1 for the sequence  $\{u_n\}$ . However, if  $\#S = 0$ , then  $\{u_n\}$  is uniformly bounded and  $\Sigma_n \rightarrow 0$ . Thus by Theorem 1,  $u_n$  coincides with the minimal solution  $\underline{u}_n$  of (1.2) for  $n$  large. It is well known that the minimal solution  $\underline{u}_n$  is stable and its Morse index is exactly 0. This contradicts to the assumption  $i_M(u_n) = 1$ , thus we have  $\#S = 1$ .  $\square$

## 2 Proof of Theorem 2

In this section, we prove Theorem 2 along the line of [13], [14]. Analytical tools needed for the study of singular Liouville equations are provided in Tarantello's nice book [12]. In the proof, we need a concentration-compactness alternative result of Bartolucci and Tarantello ([2], [3], see also [12]: Proposition 5.4.32), which we recall here in the following form.

**Proposition 4** *Let  $v_n$  satisfy*

$$\begin{aligned} -\Delta v_n &= |x-p|^{2\alpha} W_n(x) e^{v_n} \quad \text{in } B_1(p) \subset \mathbb{R}^2, \\ \int_{B_r(p)} |x-p|^{2\alpha} W_n e^{v_n} dx &\leq C \quad \text{for some } r \in (0, 1], \end{aligned}$$

where  $\alpha > 0$  and  $W_n$  is a  $C^1$  function on  $B_1(p)$  such that

$$0 < b_1 \leq W_n \leq b_2, \quad |\nabla W_n| \leq A \quad \text{in } B_1(p)$$

for some  $b_1, b_2, A > 0$  uniformly in  $n$ .

Then there exists  $\delta \in (0, 1]$  and a subsequence of  $v_n$  (denoted by the same symbol), for which only one of the following alternatives hold:

- (a)  $v_n$  is bounded uniformly in  $L_{loc}^\infty(B_\delta(p))$ ;
- (b)  $\sup_{\Omega'} v_n \rightarrow -\infty$  for every  $\Omega' \subset\subset B_\delta(p)$ ;
- (c) there exists  $z_n \in B_1(p)$  such that  $z_n \rightarrow p$  and  $v_n(z_n) \rightarrow +\infty$ , while  $\sup_{\Omega'} v_n \rightarrow -\infty$  for every  $\Omega' \subset\subset B_\delta(p) \setminus \{p\}$  and  $|x-p|^{2\alpha} W_n e^{v_n} \rightarrow \beta \delta_p$  in the sense of measures in  $B_\delta(p)$  with  $\beta \geq 4\pi$ . Furthermore if  $W_n \rightarrow W$  in  $C_{loc}^0$  for some  $W$ , then  $\beta \geq 8\pi$ .

Let  $\{u_n\}$  be a solution sequence to (1.2) for  $\lambda = \lambda_n$  with  $\Sigma_n = O(1)$  as  $n \rightarrow \infty$ . If  $\Sigma_n \rightarrow 0$ , then  $S = \emptyset$  and we have nothing to prove. Thus we consider the case (ii) of Theorem 1, and we have a blow up set  $\mathcal{S} = \{q_1, \dots, q_m\} \subset \Omega$  for (a subsequence of)  $\{u_n\}$ .

Let  $L_n = -\Delta_x - \lambda_n K(x) e^{u_n(x)} : H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$  be the linearized operator around  $u_n$  and let  $\lambda_j(L_n, D)$  denote the  $j$ -th eigenvalue of  $L_n$  acting on  $H_0^1(D)$  for a regular subdomain  $D \subset \Omega$ . Next is the key in the proof of Theorem 2.

**Claim:** There exist  $m$  disjoint open balls  $\{B^i\}_{i=1}^m$ , each  $B^i \subset\subset \Omega$ , such that  $\lambda_1(L_n, B^i) < 0$  for any  $i \in \{1, \dots, m\}$  and for  $n$  large.

Assuming for the moment the validity of Claim, we prove Theorem 2. Indeed, by Claim, there exist  $m$  open balls  $B^1, \dots, B^m$  which are disjoint, such that

$$\lambda_1(L_n, B^i) < 0 \quad \text{for } i = 1, \dots, m.$$

On the other hand, it is well known that

$$\lambda_m(L_n, \Omega) \leq \sum_{i=1}^m \lambda_1(L_n, B^i)$$

holds; see, for example, the Appendix of [13]. Combining these inequalities, we have  $\lambda_m(L_n, \Omega) < 0$ . Therefore by the definition of the Morse index of  $u_n$ , we have  $m \leq i_M(u_n)$ . This proves Theorem 2.  $\square$

In the following, we will prove Claim. Let  $S \setminus \Gamma = \{q_{i_1}, \dots, q_{i_k}\}$ . Since  $K(x) = \prod_{j=1}^N |x - p_j|^{2\alpha_j} V(x)$  is strictly positive smooth function near any  $q \in S \setminus \Gamma$ , the argument in [14], which uses a concentration-compactness result of [4] [8], works well around  $q \in S \setminus \Gamma$ . Thus we can find  $r$  disjoint balls  $\{B'_i\}_{i=1}^k$  with the desired property. We refer the reader to [14] [13].

Next, we consider blow up points in  $S \cap \Gamma = \{p_{j_1}, \dots, p_{j_s}\}$  and, for simplicity, we relabel  $S \cap \Gamma = \{p_1, \dots, p_s\}$ . We choose  $r > 0$  sufficiently small such that  $B_r(p_i) \subset\subset \Omega$ ,  $\{B_r(p_i)\}_{i=1}^s$  are disjoint, and  $p_i$  is the only blow up point of  $u_n$  in  $B_r(p_i)$  for all  $i$ . Let  $x_n^i \in B_r(p_i)$  be a point such that

$$u_n(x_n^i) = \max_{B_r(p_i)} u_n(x) \rightarrow +\infty, \quad x_n^i \rightarrow p_i \quad (i = 1, \dots, s),$$

as  $n \rightarrow \infty$ .

Now, let us define  $\delta_n^i > 0$  and  $\tilde{u}_n^i : B_{r/\delta_n^i}(0) \rightarrow \mathbb{R}$  so that

$$\begin{aligned} (\delta_n^i)^{2\alpha_i+2} \lambda_n e^{u_n(p_i)} &= 1, \\ \tilde{u}_n^i(y) &= u_n(\delta_n^i y + p_i) - u_n(p_i), \quad y \in B_{r/\delta_n^i}(0) \end{aligned}$$

for  $i \in \{1, \dots, s\}$ .

First, we prove

**Lemma 5**  $\delta_n^i \rightarrow 0$  as  $n \rightarrow \infty$ .

*Proof.* Define  $v_n(x) = u_n(x) + \log \lambda_n$ . Then  $v_n$  satisfies

$$-\Delta v_n = |x - p_i|^{2\alpha_i} \hat{K}_i(x) e^{v_n} \quad \text{in } B_r(p_i), \quad v_n = u_n + \log \lambda_n \quad \text{on } \partial B_r(p_i),$$

where  $K(x) = |x - p_i|^{2\alpha_i} \hat{K}_i(x)$ ,  $\hat{K}_i(x) = \prod_{j=1, j \neq i}^N |x - p_j|^{2\alpha_j} V(x)$ . Note that  $\hat{K}_i$  is a smooth, strictly positive function on  $B_r(p_i)$ . Also, Theorem 1 (1.3), (1.4) implies that

$$\int |x - p_i|^{2\alpha_i} \hat{K}_i(x) e^{v_n} dx \rightarrow 8\pi(1 + \alpha_i) \delta_{p_i} \quad (2.1)$$

in the sense of measures on  $B_r(p_i)$  and

$$\max_{\partial B_r(p_i)} v_n(x) - \min_{\partial B_r(p_i)} v_n(x) = \max_{\partial B_r(p_i)} u_n(x) - \min_{\partial B_r(p_i)} u_n(x) = O(1) \quad (2.2)$$

as  $n \rightarrow \infty$ . Recall the assumption  $\alpha_i \notin \mathbb{N}$  for all  $i$ . Therefore, we can apply Proposition 5.6.50 and Corollary 5.4.24 in [12] to  $v_n$  to conclude that

$$\sup_{B_\rho(p_i)} \{v_n(x) + (2\alpha_i + 1) \log |x - p_i|\} \leq C$$

for any  $\rho < r$ , which implies  $\left(\frac{|x_n^i - p_i|}{\delta_n^i}\right)^{2(\alpha_i+1)} \leq e^C$ , and

$$v_n(p_i) = \max_{B_r(p_i)} v_n + O(1) \quad (2.3)$$

as  $n \rightarrow \infty$ . Thus  $u_n(p_i) = u_n(x_n^i) + O(1) \rightarrow \infty$  for any  $i = 1, \dots, s$  as  $n \rightarrow \infty$ .

Now, we claim that  $v_n(p_i) \rightarrow +\infty$  as  $n \rightarrow \infty$  for any  $i \in \{1, \dots, s\}$ . Indeed, assume the contrary that there exists  $i \in \{1, \dots, s\}$  and a subsequence (denoted by the same symbol) such that

- (i)  $v_n(p_i) \rightarrow -\infty$ , or
- (ii)  $v_n(p_i) \rightarrow C$  for some  $C \in \mathbb{R}$ .

When (i) happens, we see by (2.3) that

$$\int_{B_r(p_i)} K(x) e^{v_n(x)} dx \leq e^{\max_{B_r(p_i)} v_n} \int_{B_r(p_i)} K(x) dx = e^{v_n(p_i) + O(1)} \int_{B_r(p_i)} K(x) dx \rightarrow 0$$

as  $n \rightarrow \infty$ . On the other hand, since  $p_i$  is the only blow up point of  $\{u_n\}$  in  $B_r(p_i)$ , (2.1) implies

$$\lim_{n \rightarrow \infty} \int_{B_r(p_i)} K(x) e^{v_n} dx \geq 8\pi(1 + \alpha_i),$$

which leads to a contradiction.

When (ii) happens, again by (2.3), we see  $\max_{B_r(p_i)} v_n = v_n(x_n^i) = O(1)$  as  $n \rightarrow \infty$ . Since  $x_n^i \rightarrow p_i$  as  $n \rightarrow \infty$ , this case can happen only when the alternative (a) in Proposition 4 occurs:  $\{v_n\}$  is bounded in  $L_{loc}^\infty(B_r(p_i))$ . On the other hand, since  $u_n = O(1)$  locally on  $B_r(p_i) \setminus \{p_i\}$  by (1.4),  $v_n = u_n + \log \lambda_n \rightarrow -\infty$  on any compact set in  $B_r(p_i) \setminus \{p_i\}$ . This again leads to a contradiction and we have proved the claim. Now, since  $(\delta_n^i)^{2(1+\alpha_i)} = \frac{1}{e^{v_n(p_i)}}$ , we obtain the lemma.  $\square$

Incidentally, by (2.1), (2.2) and (2.3), we can apply Theorem 5.6.51 in [12], see also [1], to  $v_n$  to obtain the following pointwise estimate

$$\left| v_n(x) - \log \frac{e^{v_n(p_i)}}{\left(1 + \frac{1}{8(\alpha_i+1)^2} c_i e^{v_n(p_i)} |x - p_i|^{2(\alpha_i+1)}\right)^2} \right| \leq C \quad \text{for } x \in B_r(p_i),$$

which is equivalent to

$$\left| u_n(x) - \log \frac{e^{u_n(p_i)}}{\left(1 + \frac{\lambda_n}{8(\alpha_i+1)^2} c_i e^{u_n(p_i)} |x - p_i|^{2(\alpha_i+1)}\right)^2} \right| \leq C \quad \text{for } x \in B_r(p_i),$$

where  $c_i = \hat{K}_i(p_i)$ .

Going back to the proof of Theorem 2, we see that  $\tilde{u}_n^i$  satisfies

$$\begin{cases} -\Delta \tilde{u}_n^i = |y|^{2\alpha_i} \hat{K}_i(\delta_n^i y + p_i) e^{\tilde{u}_n^i} & \text{in } B_{r/\delta_n^i}(0), \\ \hat{K}_i(\delta_n^i y + p_i) \rightarrow c_i = \hat{K}_i(p_i) & \text{uniformly in } C_{loc}^0(\mathbb{R}^2), \\ \tilde{u}_n^i(0) = 0, \max_{B_{r/\delta_n^i}(0)} \tilde{u}_n^i = u_n(x_n^i) - u_n(p_i) = O(1), \\ \int_{B_{r/\delta_n^i}(0)} |y|^{2\alpha_i} \hat{K}_i(\delta_n^i y + p_i) e^{\tilde{u}_n^i} dy = O(1), & (n \rightarrow \infty). \end{cases}$$

The third equation comes from (2.3).

At this stage, we can apply Lemma 5.4.21 in [12] to  $\tilde{u}_n^i$  to confirm that  $\tilde{u}_n^i$  is uniformly bounded in  $L_{loc}^\infty(\mathbb{R}^2)$  and along a subsequence,

$$\tilde{u}_n^i \rightarrow U^i(y) \quad \text{in } C_{loc}^2(\mathbb{R}^2) \text{ as } n \rightarrow \infty, \quad (2.4)$$

where  $U^i$  satisfies

$$\begin{cases} -\Delta U^i = c_i |y|^{2\alpha_i} e^{U^i} & \text{in } \mathbb{R}^2, \\ U^i(0) = 0, \\ \int_{\mathbb{R}^2} |y|^{2\alpha_i} e^{U^i} dy < +\infty. \end{cases}$$

By a classification result of Prajapat and Tarantello [11] and the assumption  $\alpha_i \notin \mathbb{N}$ , we have

$$U^i(y) = -2 \log \left( 1 + \frac{c_i}{8(\alpha_i+1)^2} |y|^{2(\alpha_i+1)} \right) \quad \text{for } i = 1, \dots, s.$$

Now, we define

$$\tilde{L}_n^i = -\Delta_y - |y|^{2\alpha_i} \hat{K}_i(\delta_n^i y + p_i) e^{\tilde{u}_n^i(y)} : H_0^1(B_{r/\delta_n^i}(0)) \rightarrow H^{-1}(B_{r/\delta_n^i}(0)).$$

This operator is related to  $L_n$  by the formula

$$(\delta_n^i)^2 L_n \Big|_{u_n(x) = \tilde{u}_n^i(y) + u_n(p_i)} = \tilde{L}_n^i,$$

where  $x = \delta_n^i y + p_i$  for  $x \in B_r(p_i)$  and  $y \in B_{r/\delta_n^i}(0)$ . Also for a domain  $D \subset B_r(p_i)$ , we have

$$(\delta_n^i)^2 \lambda_1(L_n, D) = \lambda_1(\tilde{L}_n^i, D_n^i), \quad D_n^i = \frac{D - p_i}{\delta_n^i}, \quad (2.5)$$

where  $\lambda_1(\tilde{L}_n^i, D_n^i)$  denotes the first eigenvalue of  $\tilde{L}_n^i$  acting on  $H_0^1(D_n^i)$ .

Now, we show

**Lemma 6** *There exist disjoint balls  $\{B_{\delta_n^i R}(p_i)\}_{i=1, \dots, s}$  for some  $R > 0$  such that  $\lambda_1(L_n, B_{\delta_n^i R}(p_i)) < 0$  for  $n$  large and for any  $i \in \{1, \dots, s\}$ .*

**Proof.** For  $R > 0$ , we define

$$w_R(y) = 2 \log \frac{8 + R^2}{8 + |y|^2} \in H_0^1(B_R(0)).$$

We will prove that  $(\tilde{L}_n^i w_R, w_R)_{L^2(B_R)} < 0$  for  $n \in \mathbb{N}$  and  $R > 0$  sufficiently large with  $B_R(0) \subset B_{r/\delta_n^i}(0)$ . Indeed,

$$\begin{aligned} (\tilde{L}_n^i w_R, w_R)_{L^2(B_R)} &= \int_{B_R(0)} |\nabla w_R|^2 dy - \int_{B_R(0)} |y|^{2\alpha_i} \hat{K}_i(\delta_n^i y + p_i) e^{\tilde{u}_n^i(y)} w_R^2(y) dy \\ &=: I_1 - I_2. \end{aligned}$$

We observe that

$$I_1 = \int_{B_R(0)} \frac{16|y|^2}{(8 + |y|^2)^2} dy = 2\pi \int_0^R \frac{16r^2}{(8 + r^2)^2} r dr \leq 32\pi (\log R) [1 + o_R(1)],$$

where  $o_R(1) \rightarrow 0$  as  $R \rightarrow \infty$ . On the other hand, we have

$$\begin{aligned} I_2 &= \int_{B_R(0)} |y|^{2\alpha_i} \hat{K}_i(\delta_n^i y + x_n^i) e^{\tilde{u}_n^i(y)} w_R^2(y) dy \\ &= c_i \int_{B_R(0)} \frac{|y|^{2\alpha_i}}{\left(1 + \frac{c_i}{8(\alpha_i+1)^2} |y|^{2\alpha_i+2}\right)^2} \left\{2 \log \frac{8 + R^2}{8 + |y|^2}\right\}^2 dy + o_n(1) \\ &= 8\pi c_i \int_0^R \frac{r^{2\alpha_i+1}}{\left(1 + \frac{c_i}{8(\alpha_i+1)^2} r^{2\alpha_i+2}\right)^2} \left\{\log(8 + R^2) - \log(8 + r^2)\right\}^2 dr + o_n(1) \\ &= 8\pi c_i \left[ \frac{4(\alpha_i + 1)}{c_i} + o_R(1) \right] \left\{\log(8 + R^2)\right\}^2 + o_n(1) \\ &= 32\pi(\alpha_i + 1) \left\{\log(8 + R^2)\right\}^2 [1 + o_R(1)] + o_n(1), \end{aligned}$$

where we have used (2.4) and

$$\int_0^R \frac{r^{2\alpha+1}}{(1+cr^{2\alpha+2})^2} dr = \int_0^\infty \frac{r^{2\alpha+1}}{(1+cr^{2\alpha+2})^2} dr + o_R(1) = \frac{1}{2(\alpha+1)c} + o_R(1)$$

for  $\alpha, c > 0$ . Thus we obtain

$$(\tilde{L}_n^i w_R, w_R)_{L^2(B_R)} = I_1 - I_2 \leq -32\pi(\alpha_i + 1) \{\log(8 + R^2)\}^2 [1 + o_R(1)] < 0$$

by taking  $n$  sufficiently large first, and then  $R > 0$  large such that  $B_R(0) \subset B_{r/\delta_n^i}(0)$ . This implies that the first eigenvalue of the operator  $\tilde{L}_n^i$  on  $B_R$  is negative:  $\lambda_1(\tilde{L}_n^i, B_R) < 0$ . By this calculation and (2.5) proves that  $\lambda_1(L_n, B_{\delta_n^i R}(p_i)) < 0$  for  $i = 1, \dots, s$ . These balls  $\{B_i\}_{i=1}^s = \{B_{\delta_n^i R}(p_i)\}_{i=1}^s$  can be disjoint if we choose sufficiently large  $n$ , since the blow up set  $S$  is finite and  $\delta_n^i = o(1)$  as  $n \rightarrow \infty$ .  $\square$

Since balls  $\{B_i\}_{i=1}^s$  in Lemma 6 can also be made disjoint from balls  $\{B'_l\}_{l=1}^k$  (former obtained around points in  $S \setminus \Gamma$ ), we obtain Claim. The proof of Theorem 2 is completed.  $\square$

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