

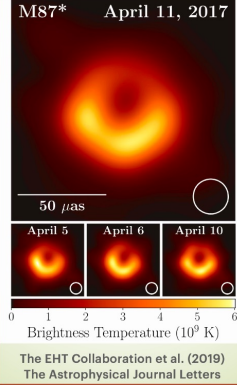
Gravitational Lensing in a Perfect Fluid Plasma Free-Falling in the Kerr Spacetime

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1. Introduction & Background

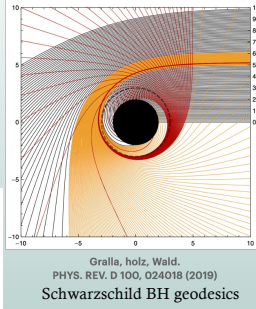
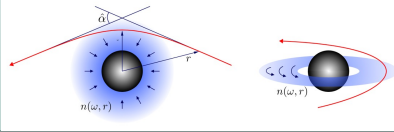
- EHT project observes black hole candidates like M87* and Sgr A*
- Since electromagnetic waves traverse a plasma surrounding the black hole candidate, the rays are affected by the plasma
- These effects are frequency-dependent = chromatic



✓ Researching astrophysical objects surrounded by plasma is important both theoretically and observationally.

2. Research Goals

1. **Goal:** to compute light trajectories around the Kerr BH with plasma density that gives a **non-separable equation**.
2. The plasma four-velocity is considered including its radial and azimuthal motion [1]. Please imagine a rainbow.



3. Setup (Kerr Black Hole)

- Gravitational lensing + dispersive medium (plasma)

The Kerr metric can represent a spacetime around a rotating and uncharged black hole (Asymptotically flat, axisymmetric, stationary).

$$g_K = g_{tt}dt^2 + 2g_{t\phi}dtd\phi + g_{rr}dr^2 + g_{\phi\phi}d\phi^2 + g_{\theta\theta}d\theta^2$$

$$\begin{aligned} \Sigma &= \rho^2 = r^2 + a^2 \cos^2 \theta, \\ \Delta &= r^2 - 2Mr + a^2, \\ A &= (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta, \\ g_{tt} &= -\left(1 - \frac{2Mr}{\Sigma}\right), \quad g_{t\phi} = -2Mar \sin^2 \theta / \Sigma, \\ g_{rr} &= \Sigma / \Delta, \quad g_{\theta\theta} = \Sigma, \\ g_{\phi\phi} &= A \sin^2 \theta / \Sigma, \end{aligned}$$

$$\hat{h} = \begin{pmatrix} g_{tt} & g_{t\phi} \\ g_{\phi t} & g_{\phi\phi} \end{pmatrix} \quad x = r^2 + a^2$$

- The medium is specified by its refractive index n (which is the reciprocal of the phase speed) and the medium's four-velocity V^i . Labels = 0, 1, 2, 3
- The canonical conjugate momenta $p_\mu = \frac{\partial L}{\partial v^\mu}$ or $p_\mu = \frac{\partial S}{\partial x^\mu}$
- The refractive index is a given function of coordinates x_i and frequency ω : $n \equiv n(x^i, \omega) \quad \omega \equiv \omega(p_i, x^i) = -p_i V^i$
- Using the relation $n^2 = 1 + \frac{p_i p^i}{(p_k V^k)^2}$, and the ray condition $\mathcal{H}(x^i, p_i) = 0$, the Hamiltonian (Synge's text [2]) becomes
$$\mathcal{H}(x^i, p_i) = \frac{1}{2} \{ g^{ik} p_i p_k - [n^2(x^i, \omega(p_i, x^i)) - 1] (p_j V^j)^2 \}$$

4. Perfect fluid around the Kerr BH

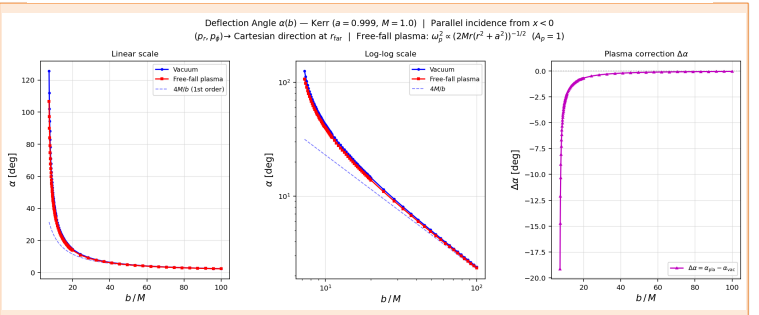
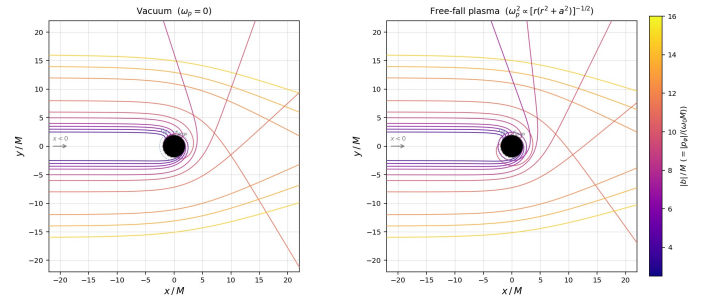
- ✓ The Hamiltonian of geodesics: $H = \frac{1}{2} g^{ab} p_a p_b$
- ✓ Carter (1968) proposed the separable action of the Hamilton-Jacobi equation for geodesics in the Kerr spacetime. Let μ be the rest mass, E be the energy, L_z be the angular momentum along the symmetry axis, and Q be the Carter constant.
- ✓ We impose the conditions $E = \mu = 1, L_z = Q = 0$ on the constants, which is the free fall from infinity.

- ✓ The number density $N(r, \theta)$ and the four-velocity U^a should be stationary and axisymmetric. Solve the continuity equation $\nabla_a (N U^a) = 0$, which provides the plasma frequency
$$\omega_p^2 = \frac{e^2}{\epsilon_0 m_e} N(r, \theta),$$
- ✓ Tildes mean non-dimensional
$$\tilde{\omega}_p^2 \equiv \frac{\omega_p^2}{\omega_0^2} = \frac{A(\theta)}{\sqrt{2\tilde{r}\tilde{x}}}$$
- ✓ $A(\theta)$ is a (constant) function.
- ✓ $\omega_0 = -p_t, N \propto 1/\sqrt{r(r^2 + a^2)}$
- ✓ Perlick and Tsupko (2017) derived the separability condition of light trajectories but our plasma distribution does not satisfy it [3, 4].

5. Numerical calculations

Ray Tracing in Kerr Spacetime ($a = 0.5M$) — Parallel incidence from $x < 0$ | Free-fall plasma: $n = [r(r^2 + a^2)]^{-1/2}$

Runge-Kutta method (Python) or Tsitouras (Julia lang.) method



6. Summary and Future work

- We computed gravitational lensing in Kerr spacetime filled with a **free-falling** plasma, using Synge's Hamiltonian in the geometrical-optics limit.
- **The plasma reduces the deflection angle** relative to vacuum at a fixed photon frequency ω_0 .
- Outlook: Circular orbits? Frequency band? Accretion disk?